

Non Linear stability of Equilibrium points in the Generalized photogravitational Restricted Three Body Problems through Kolmogonov -Arnold Moser Theory

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Abstract- *In this Paper we are presented a key points of criticism of KAM (Kolmogorov-Arnold- Moser) theory case of Generalized photo gravitational restricted three-body problem when small primaries is oblate spheroid and bigger primaries is radiated.we should especially note that there is no analogue of Jacobian-type integral of motion in the case of photo gravitational restricted three-bodt problem if we take into consideration small Yarkovsky e ect. The KAM-theory is that an appropriate Hamilton formalism should be valid for the KAM thoery , but Hamilton formalism could not be applied for restricted three-body problem (which is proved to have only the Jacobian type integral of motion, but the integrals of energy, momentum are not invariants).*

Keywords- KAM theory, Hamilton formalism, Yarkovsky e ect, photo gravitational restricted three-body problem, Jacobian - type integral of motion

I. INTRODUCTION

KAM theory is known to be applied for research of stability of solar system in term of R3BP If we consider R3BP with additional in uence of Yarkovsky e ect of non-gravitational nature on equation of motion. The Yarkovsky e ect is a force acting on a rotating body in space caused by the an isotropic emission of thermal photons, which carry momentum. It is usually considered in relation to meteoroids or small asteroids as its in uence is most signi cant for these bodies. Such a force is produced by the way an asteroid absorbs energy from the sun and re-radiates it into space as heat by an isotropic way. In fact, there exists a disbalance of momentum when asteroid at rst absorbs the light, radiating from the sun, but then asteroid re-radiates the heat. Such a disbalance is caused by the rotating of asteroid during period of warming as well as it is caused by the isotropic cooling of surface inner layers; the processes above depend on an isotropic heat transfer in the inner layers of asteroid. Thus,Yarkovsky e ect is small but very important e ect in celestial mechanics as well as in calculating of a proper orbits of asteroids other small bodies. Besides, Yarkovsky e ect is not predictable (it could be only observed measured by astronomical methods); the main reason

is un-predictable character of the rotating of small bodies in the case when there is no any collision between them.

Modi ed Equations of Motion (YORP-E ect)

Modi ed equations of motion for the three dimensional restricted 3-bodies problem, with an oblate primary m₂, both primaries radiating, and the in-nitesimal mass m under the in uence of YORP-e ect, should be presented in barycentric rotating co-ordinate system in the form below:

These equations have the form

$$\ddot{x} - 2\dot{y} = \frac{\partial}{\partial x} Yx(t) \quad \text{---}$$

$$\ddot{y} + 2\dot{x} = \frac{\partial}{\partial y} Yy(t) \quad \text{---}$$

$$\ddot{z} = \frac{\partial}{\partial z} Yz(t)$$

where $Yx(t)$, $Yy(t)$, $Yz(t)$ are the projecting of YORP-e ect acceleration $Y(t)$ on the appropriate axis Ox , Oy , Oz .

Location of Equilibrium Points

The location of equilibrium points for system (in general is given by con-ditions $\dot{x} = \dot{y} = \dot{z} = x = y = z = 0$

$$\frac{\partial}{\partial x} = - \underline{Yx(t)}$$

$$\frac{\partial}{\partial y} = - \underline{Yy(t)}$$

$$\frac{\partial}{\partial z} = - \underline{Yz(t)}$$

which determine the location of equilibrium points

we should note that a case of Yarkovsky e ect is negligible determines the existence of quasi-planar

equilibrium points in which conditions Yz approaches to 0, z approaches to 0 are valid simultaneously. The strongest simplifying of system is possible when Yarkovsky effect is zero, $Yx = Yy = Yz = 0$. In such a case, it has been proved the existence equilibrium points in photogravitational restricted 3-bodies problem.

STABILITY

Now we apply Mosers modified form of Arnolds theorem to discuss the non linear stability that equations of restricted three-body problem are proved to describe the system with non-Hamilton formalism

$$= \frac{x^2+y^2}{2} + \frac{1}{r_1} + \frac{1}{r_2}$$

These equations have the form

$$\dot{x} - 2y = \frac{\partial U}{\partial x}$$

$$\dot{y} + 2x = \frac{\partial U}{\partial y}$$

$$\dot{z} = \frac{\partial U}{\partial z}$$

For example

$$\frac{\partial U}{\partial x} = X - \frac{(1-\mu)r_1}{r_1^3} - \frac{\mu r_2}{r_2^2}$$

where

$$r_1^2 = (x + \mu)^2 + y^2 + z^2$$

$$2r_1 \frac{\partial r_1}{\partial x} = 2(x + \mu)$$

$$\frac{\partial r_1}{\partial x} = \frac{x + \mu}{r_1}$$

$$r_2^2 = (x - 1 + \mu)^2 + y^2 + z^2$$

$$\frac{\partial r_2}{\partial x} = \frac{x - 1 + \mu}{r_2}$$

Substituting Equations and into Equation yields the first of equations

$$\dot{x} - 2y = X - (1 - \mu) \frac{(x + \mu)}{r_1^2} - \frac{\mu(x - 1 + \mu)}{r_2^2}$$

By multiplying Equations $2x$, $2y$ and $2z$, respectively

$$2x\dot{x} - 22xy = 2x \frac{\partial U}{\partial x}$$

$$2y\dot{y} + 22xy = 2y \frac{\partial U}{\partial y}$$

$$\frac{\partial U}{\partial y} 2z\dot{z} = 2z \frac{\partial U}{\partial z}$$

Summing these

$$2x\dot{x} + 2y\dot{y} + 2z\dot{z} = 2x \frac{\partial U}{\partial x} + 2y \frac{\partial U}{\partial y} + 2z \frac{\partial U}{\partial z}$$

This can be integrated to yield

$$x^2 + y^2 + z^2 = 2C$$

We have used the fact that

$$2 \left[\frac{dx}{dt} \frac{\partial U}{\partial x} + \frac{dy}{dt} \frac{\partial U}{\partial y} + \frac{dz}{dt} \frac{\partial U}{\partial z} \right]$$

We represent a diffusion mechanism for the time - dependent perturbation of autonomous Hamiltonian system introduced unlike other approach does not rely on the KAM theory. According to above constant, we could obtain from the equations of system a Jacobian-type integral of motion

Besides, the restricted three-body problem is proved to have a new, the only stable invariant = Jacobian-type integral of motion.

Hamilton systems are assumed to be the systems without diffusion. It means that such a dynamical systems should have a weak Arnold-diffusion.

- where C is so-called Jacobian constant. As it was proved such a Jacobian type integral of motion should not be

depending on time for large time-period. Additionally, we should especially note obvious fact: in the case of photogravitational restricted three-body problem with Yarkovsky effect there is no analogue of Jacobian-type integral for ODE system of motion.

II. CONCLUSION

KAM theory tried to predict the stability in generalization photogravitational restricted three body problem. We consider solar system as Arnold division there is no analogue when we consider very small Yarkovsky effect as we discuss above that stable invariant behave as Jacobian constant. It proved that the existence of liberation points in R3BP when we take into consider even a small Yarkovsky effect.

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